

DIM-ACAV COLLOQUIUM

Cosmology with weak-lensing peak counts



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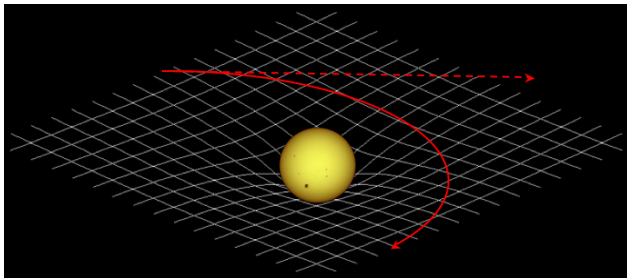


Observatoire de Paris
December 2nd, 2015

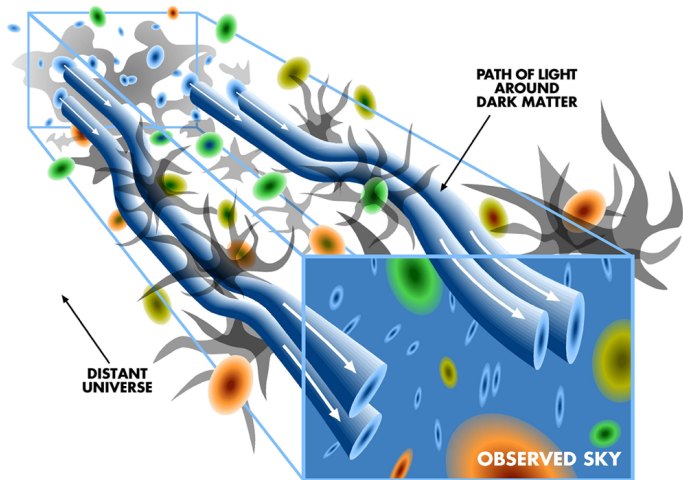
This work is supported by grants from Région Île-de-France.

Light deflection

General relativity

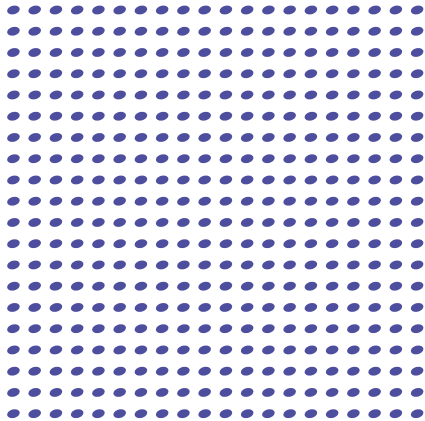


Weak lensing

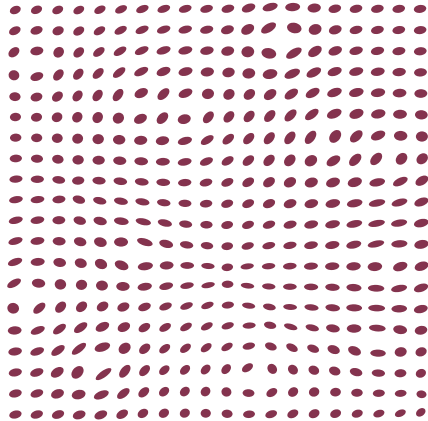


(Source: LSST)

Weak lensing



Unlensed sources



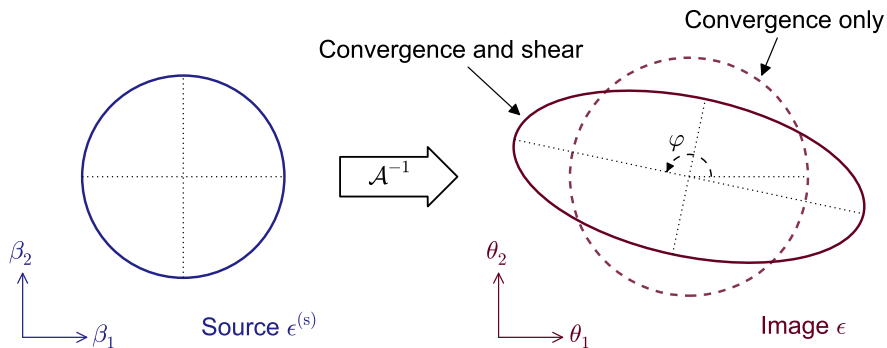
Weak lensing

Lensing formalisms

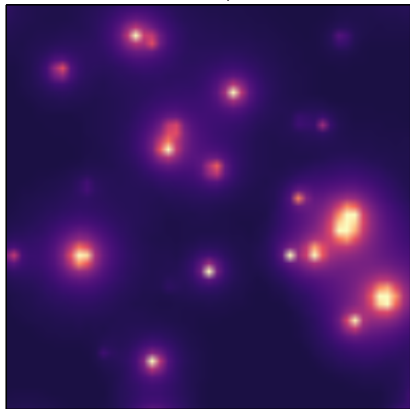
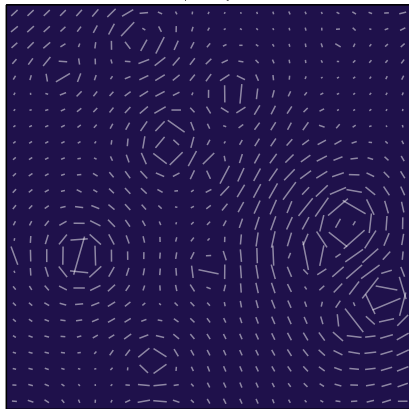
$$\mathcal{A}(\boldsymbol{\theta}) = \begin{pmatrix} 1 - \kappa - \gamma_1 & -\gamma_2 \\ -\gamma_2 & 1 - \kappa + \gamma_1 \end{pmatrix}$$

κ : convergence, “projected mass density”

$\boldsymbol{\gamma} = \gamma_1 + i\gamma_2$: cosmic shear, distortion

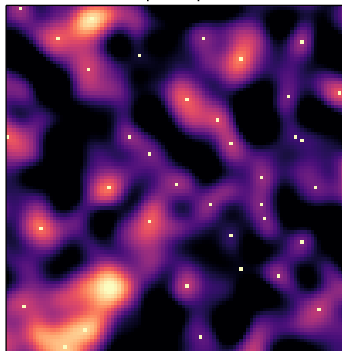


Convergence and shear maps

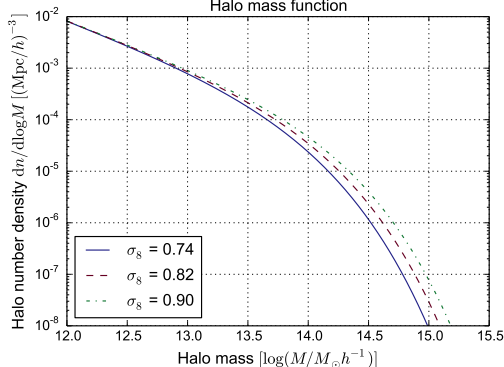
 κ map γ map

What are peak counts?

κ map and peaks



Halo mass function



- Local maxima of the projected mass
- Probe the mass function
- Contain non-Gaussian information

Problematic of peak-count modelling

Analytical models for WL peaks

- Maturi et al. (2010) [0907.1849] and Fan et al. (2010) [1006.5121]
- Difficult to handle realistic scenarios: masking, photo- z errors, etc.
- Fail to include additional features: baryons, intrinsic alignment, etc.
- Still require external simulations for covariances

Modelling with N -body simulations

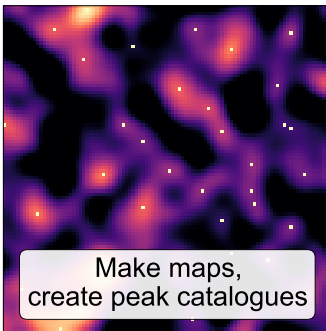
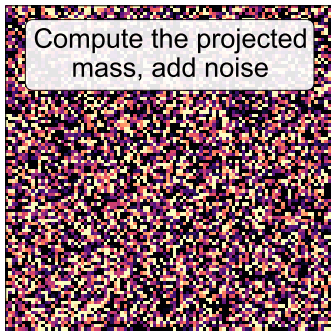
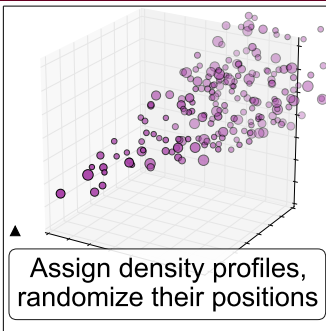
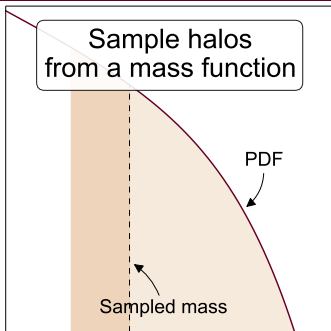
- Very expensive time costs

Challenges

How to model weak-lensing peak counts properly in realistic conditions?

How to extract cosmological information from peaks?

A new model to predict WL peak counts



A new model to predict WL peak counts

Hypotheses

- Diffuse, unbound matter does not significantly contribute to peak counts
- Spatial correlation of halos has a minor influence

Public code



CAMELUS

Counts of Amplified Mass Elevations
from Lensing with Ultrafast Simulation

<http://www.cosmostat.org/software/camelus/>

Advantages

Fast

Flexible

Full PDF information

Advantages

Fast

Only few seconds for creating a 25-deg² field, without MPI or GPU programming

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Full PDF information

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Flexible

Straightforward to include observational effects (e.g. photo- z errors, masks) and additional features (e.g. intrinsic alignment, baryonic effects)

Full PDF information

Advantages

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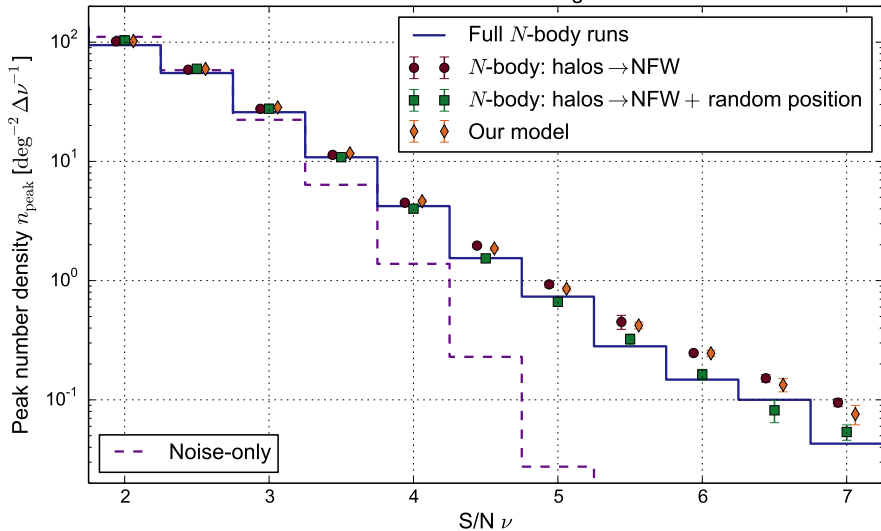
Straightforward to include observational effects (e.g. photo- z errors, masks) and additional features (e.g. intrinsic alignment, baryonic effects)

Full PDF information

Free from the Gaussian likelihood assumption, allow more flexible constraint methods (copula, varying covariances, p -value evaluation, approximate Bayesian computation, etc.)

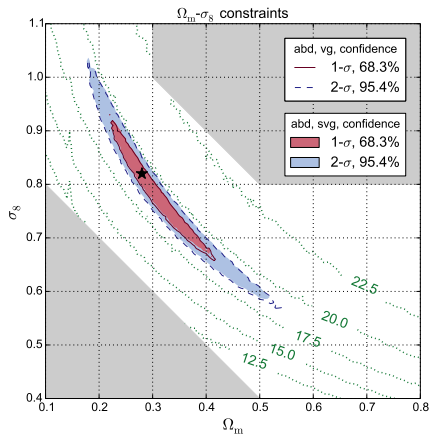
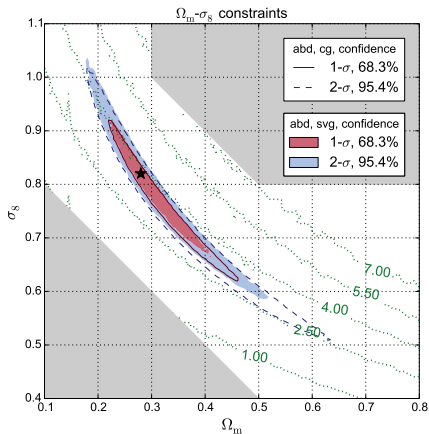
Validation

Peak abundance histogram



Lin & Kilbinger (2015a)

Cosmology-dependent covariances



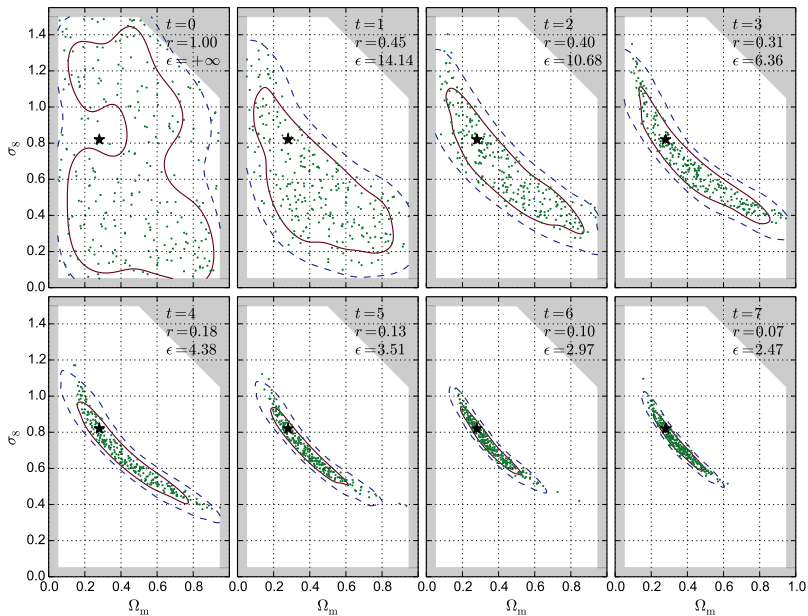
cg = constant covariance (lines, left panel)
 svg = semi-varying covariance (colored zones)
 vg = varying covariance (lines, right panel)

| | cg | svg | vg |
|-----|----|-----|----|
| FoM | 46 | 57 | 56 |

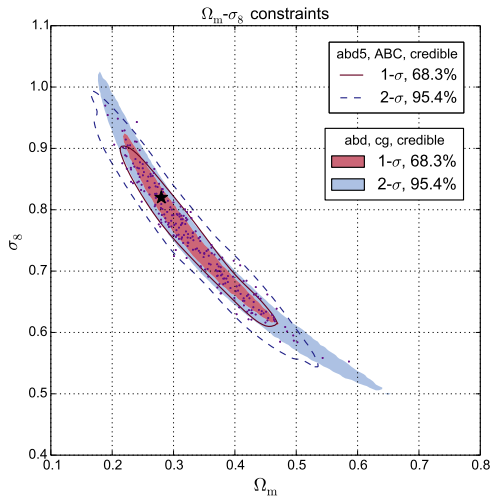
Lin & Kilbinger (2015b)

Approximate Bayesian computation

PMC ABC posterior evolution



Approximate Bayesian computation



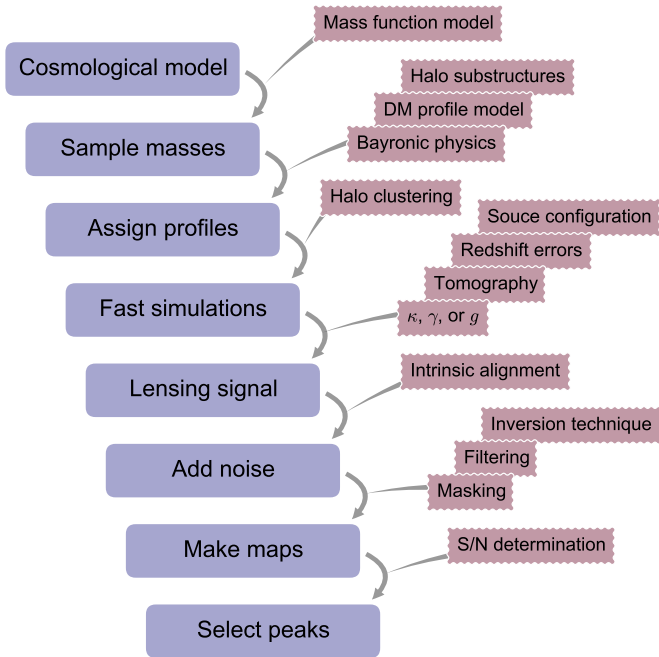
Very good agreement

Time comparison:

- Likelihood $\approx 8000 \times 1000$ simulations to run
- ABC $\approx 250 \times 1 \times 100$ simulations to run

Lin & Kilbinger (2015b)

Perspectives

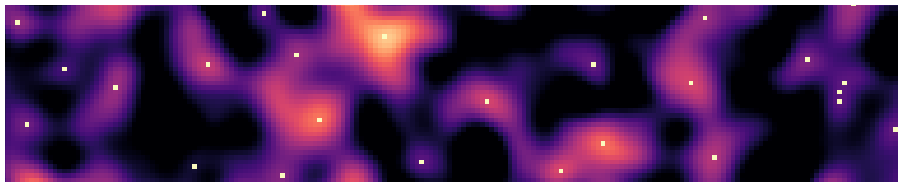


In the framework of this thesis

- Validation of the model (Lin & Kilbinger 2015a)
- Parameter constraint strategies (Lin & Kilbinger 2015b)
- Nonlinear filtering (Lin et al. 2016 in prep.)
- Application to CFHTLenS and KiDS data

Summary

Lin & Kilbinger (2015a,b) — [1410.6955] & [1506.01076]



A new model to predict WL peak counts:

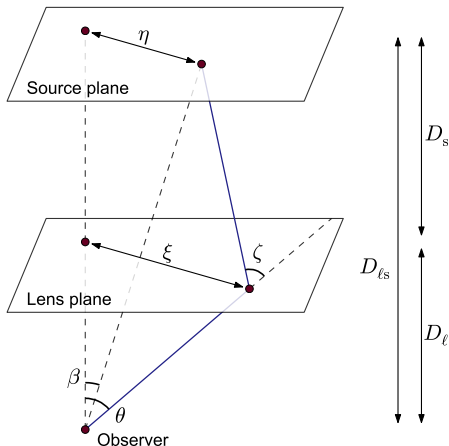
Fast, Flexible, Full PDF information

A robust and efficient constraining method:

Approximate Bayesian computation

Backup slides

Born approximation



$$A_{ij}(\boldsymbol{\theta}) = \delta_{ij} - \frac{2}{c^2} \int_0^w dw' \frac{f_K(w-w')f_K(w')}{f_K(w)} \cdot \phi_{,ij}(f_K(w')\boldsymbol{\theta}, w')$$

$$\kappa(\boldsymbol{\theta}, w) = \frac{3H_0^2\Omega_m}{2c^2} \int_0^w dw' \frac{f_K(w-w')f_K(w')}{f_K(w)} \cdot \frac{\delta(f_K(w')\boldsymbol{\theta}, w')}{a(w')}$$

Settings

- Fixed source redshift $z_s = 1.0$
- Galaxy number density $n_g = 25 \text{ arcmin}^{-2}$
- Pixel size $\theta_{\text{pix}} = 0.2 \text{ arcmin}$
- Uncorrelated Gaussian shape noise, no IA
- No mask, no baryon
- Gaussian smoothing, with radius $\theta_G = 1 \text{ arcmin}$

Validation

We compare the following four cases:

Case 1: full N -body runs

Case 2: replace N -body halos with NFW profiles with the same mass

Case 3: randomize angular positions of halos from Case 2

Case 4: our model

to test our two hypotheses:

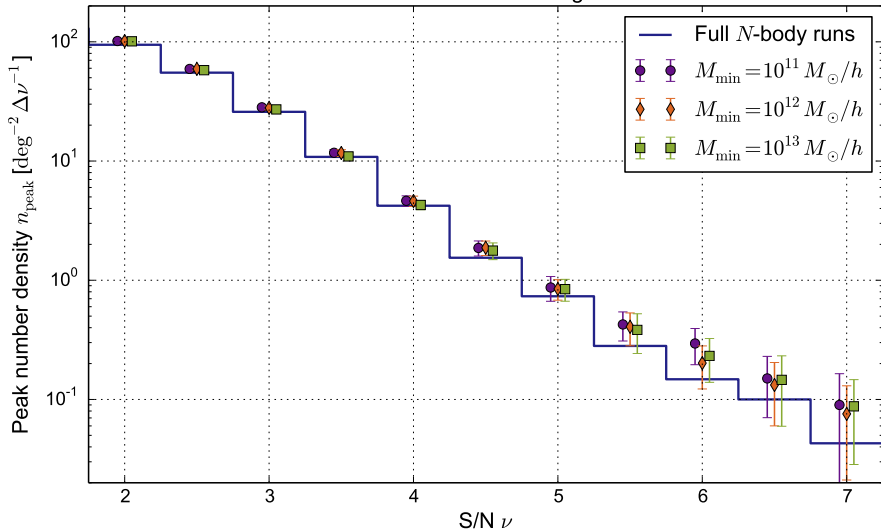
Comparison 1 & 2: contribution of unbound matters & halo asphericity

Comparison 2 & 3: impact of the spatial correlation

Comparison 3 & 4: mass function

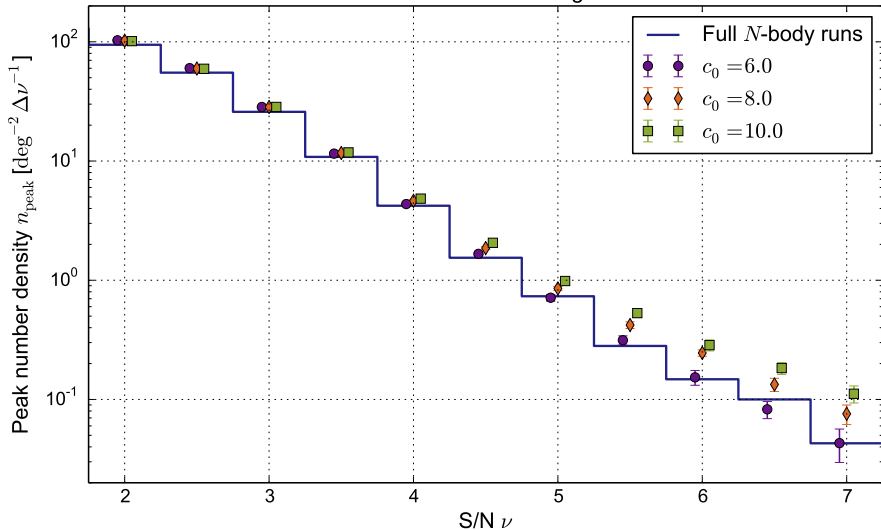
Low-mass halos seem to be negligible

Peak abundance histogram

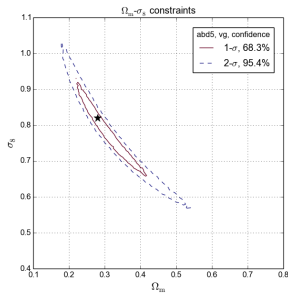


Halo structure parameters should be included in constraints

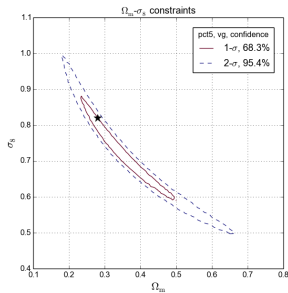
Peak abundance histogram



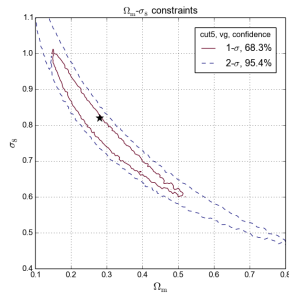
Choice of data vector



Histogram



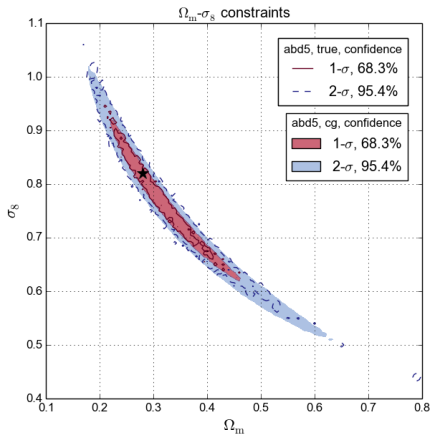
Percentile values



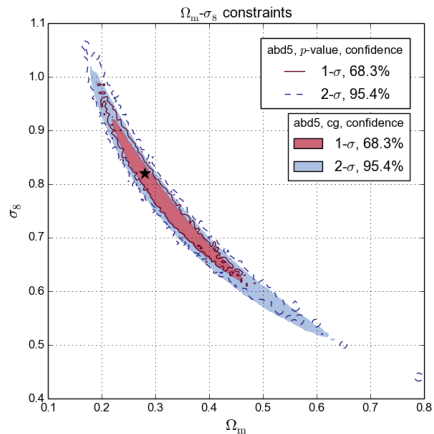
Percentile with cutoff

Lin & Kilbinger (2015b)

Non-parametric methods



With the true likelihood



With p -value

Lin & Kilbinger (2015b)

Approximate Bayesian computation

The accept-reject process of ABC samples under $\mathcal{P}_\epsilon(\pi|x^{\text{obs}})$ (red curve),

$$\mathcal{P}_\epsilon(\pi|x^{\text{obs}}) = A_\epsilon(\pi)\mathcal{P}(\pi),$$

where $P(\pi)$ is the prior (blue curve) and

$$A_\epsilon(\pi) \equiv \int dx P(x|\pi) \mathbb{1}_{|x-x^{\text{obs}}|\leq\epsilon}(x),$$

is the accept probability under π (green area).

Meanwhile,

$$\lim_{\epsilon \rightarrow 0} A_\epsilon(\pi_0)/\epsilon = P(x^{\text{obs}}|\pi_0) = \mathcal{L}(\pi_0),$$

so \mathcal{P}_ϵ is proportional to the true posterior when $\epsilon \rightarrow 0$.

